

# Contact Engineering in 2D-Material-Based Electrical Contacts

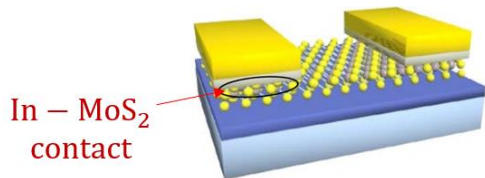
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## Motivation and Background

- The undesirably large contact resistance between two-dimensional (2D) semiconductor and three-dimensional (3D) metallic electrodes represents one of the major obstacles towards the development of practical 2D electronic and optoelectronic devices [1].
- Extensive efforts have been made to improve current flow through contacts in 2D material-based devices [2,3,4,5].
- We develop a model that includes both, the material properties of 2D semiconductors and the geometrical electrostatic effect in mixed-dimensional nanostructures, which remains rarely studied in the literature so far.

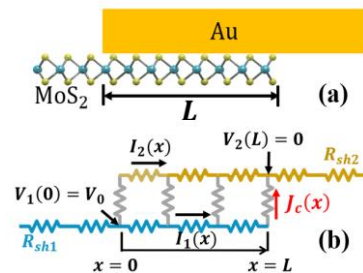


**Fig. 1:** Bottom-gate FET. The electrodes consist of 10-nm-thick In capped with 100-nm-thick Au. The ellipse under the contact indicates the interface region [5].

- [1] Allain et al., *Nat Mater* **14**, 1195 (2015)  
 [2] Das et al., *Nano Lett.* **13**, 100 (2013)  
 [3] Xu et al., *ACS Nano* **10**, 4895 (2016)  
 [4] Liu et al., *Nature* **557**, 696 (2018)  
 [5] Wang et al., *Nature* **568**, 70 (2019)

## Method

Self-consistent solutions are obtained by solving modified transmission line model [6] and thermionic emission model for 2D materials [6,7].



**Fig. 2:** (a) A typical parallel contact between Au and MoS<sub>2</sub>. (b) its transmission line model. The width (transverse dimension) of the two contact members is  $w$ .

Applying Kirchoff's law in Fig. 2b,

$$\frac{\partial V_1(x)}{\partial x} = -\frac{I_1(x)R_{sh1}}{w}, \quad \frac{\partial V_2(x)}{\partial x} = -\frac{I_2(x)R_{sh2}}{w},$$

$$\frac{\partial I_1(x)}{\partial x} = -wJ_c(x), \quad \frac{\partial I_2(x)}{\partial x} = wJ_c(x) \quad [6]$$

$$J_c(x) = J_{th} \left[ \exp\left(\frac{eV}{k_B T}\right) - 1 \right], \quad V = V_1 - V_2$$

$$J_{th} = \frac{2e\Phi_{B0}k_B T}{\pi\tau\hbar^2 v_F^2} \left( 1 + \frac{k_B T}{\Phi_{B0}} \right) \exp\left(-\frac{\Phi_{B0} - \epsilon_F}{k_B T}\right) \quad [7]$$

$$I_1(x) + I_2(x) = \text{constant} = I_{tot}$$

$$\text{Contact resistance } R_c = \frac{V_1(0) - V_2(L)}{I_{tot}}$$

$\Phi_{B0}$  = Intrinsic Schottky Barrier Height (SBH)

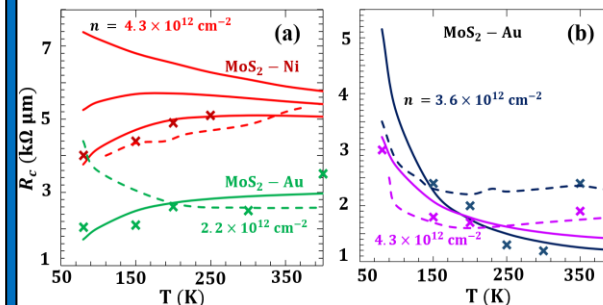
$\tau$  = carrier injection time

$v_F$  = Fermi Velocity

$\epsilon_F$  = Fermi energy

- [6] Banerjee et al., *Sci. Rep.*, **9**, 14484 (2019)  
 [7] Ang et al., *Phys. Rev. Lett.*, **121**, 056802 (2018)

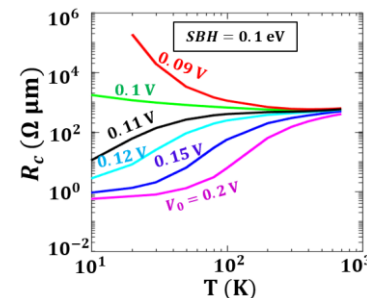
## Validation [8]



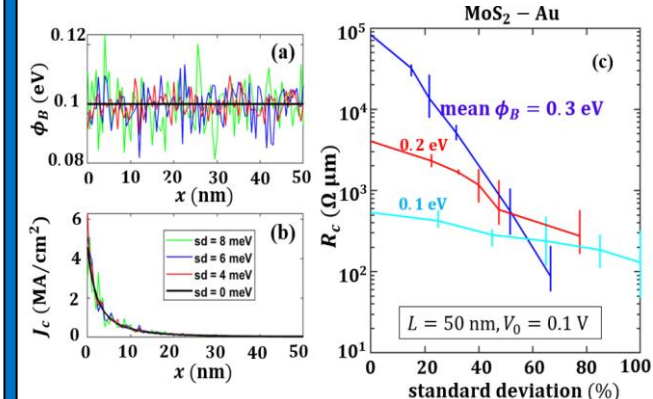
**Fig. 3:**  $R_c$  as a function of temperature  $T$  with (a) an increasing trend, and (b) a decreasing trend, for MoS<sub>2</sub> – metal contacts with different  $n$ . Crossed symbols are from experiments [9]; solid lines are from our self-consistent Model, and dashed lines are extracted from model calculations in Ref. [9]

[9] English et al., *Nano Lett.* **16**, 3824 (2016)

## Results [8]



**Fig. 4:**  $R_c$  as a function of temperature  $T$  for different applied voltage  $V_0$ , for Schottky Barrier Height (SBH) = 0.1 eV. Here, the contact is between a monolayer MoS<sub>2</sub> and Au, with  $\tau$  = 0.1 ps, and  $L$  = 50 nm.



**Fig. 5:** (a) Roughness in Schottky barrier height  $\phi_B$ , the resulting (b) current density  $J_c(x)$ , across the contact interface for a monolayer MoS<sub>2</sub>-Au 2D/3D contact for different standard deviations (sd). (c) Contact resistance  $R_c$  as a function of surface roughness (standard deviation/ $\phi_B$ ) for different mean values of  $\phi_B$ . Here,  $V_0$  = 0.1 V, and  $L$  = 50 nm.

## Conclusions [8]

- A self-consistent model to quantify the current distribution and contact resistance in 2D materials-based contacts is constructed and validated with existing experimental works.
- It is found that interface roughness can significantly reduce 2D/3D electrical contact resistance.

[8] Banerjee et al., *Phys. Rev. Appl.* **13**, 064021 (2020)

## Acknowledgement

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